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Search for Light Pseudoscalar Bosons, Pair-Produced in Higgs Boson Decays in the Four-Electron Final State in Proton-Proton Collisions at $\sqrt{s} = 13$ TeV

A. Hayrapetyan *et al.*[*]
(CMS Collaboration)

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A search for pairs of light neutral pseudoscalar bosons (A) resulting from the decay of a Higgs boson is performed. The search is conducted using LHC proton-proton collision data at $\sqrt{s} = 13$ TeV, collected with the CMS detector in 2016–2018 and corresponding to an integrated luminosity of 138 fb^{-1}. The A boson decays into a highly collimated electron-positron pair. A novel multivariate algorithm using tracks and calorimeter information is developed to identify these distinctive signatures, and events are selected with two such merged electron-positron pairs. No significant excess above the standard model background predictions is observed. Upper limits on the branching fraction for $H \to AA \to 4e$ are set at 95% confidence level, for masses between 10 and 100 MeV and proper decay lengths below 100 μm, reaching branching fraction sensitivities as low as 10^{-5}. This is the first search for Higgs boson decays to four electrons via light pseudoscalars at the LHC. It significantly improves the experimental sensitivity to axionlike particles with masses below 100 MeV .

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The standard model (SM) of particle physics has been very successful at describing the fundamental particles and their interactions. However, many questions remain unanswered, such as the strong CP problem [1,2] and the nature of dark matter (DM) [3–5]. Particles and interactions beyond the SM are postulated to solve the above-mentioned limitations. For example, an elegant solution to the strong CP problem is provided by the Peccei-Quinn mechanism [2,6,7], which introduces the axion. While the mass of the axion had originally been thought to be vanishingly small, well below 1 eV, recent studies indicate that mass values for the axion in the tens-of-MeV range are theoretically viable [8,9]. In certain DM models, the pseudoscalar axionlike particles (ALPs) [2,6,7,10] are considered as mediators of the interaction between DM and SM particles, or even as direct candidates for DM particles, especially when they are light and weakly coupled to SM particles.

Consequently, searches for axions and light ALPs (collectively referred to as ALPs, hereafter) are of great interest and have been carried out extensively [11]. Furthermore, Ref. [8] examines the implications of ALPs with a mass around 10 MeV, and hypothesizes an axion with a mass of 17 MeV from an observed excess in the electron-positron pair invariant mass spectrum, reported by the ATOMKI Collaboration [12] and recently examined by the MEG II [13] and PADME [14] experiments. Finally, it is notable that the CERN LHC has a unique potential for the observation of ALPs, especially if the Higgs boson couples to them and decays into them [15,16].

For sub-GeV particles, electromagnetic decays can be used to both trigger and cleanly identify the events in LHC experiments. There are two possible electromagnetic decay channels for ALPs: electron-positron pairs and photon pairs. In the sub-GeV mass range, ALPs can be long lived because of their small mass and decay width. LHC searches from the CMS [17] and ATLAS [18] Collaborations have set the only limits on ALPs originating in Higgs boson decays with masses down to $\mathcal{O}(100)$ MeV. Searches for ALPs coming from Higgs boson decays, with masses as low as 10 MeV, or for ALPs decaying into electron-positron pairs, are unexplored. The reason is that the decay products of such low-mass ALPs are highly collimated for typical ALP momenta above 30 GeV. The search, therefore, needs innovative techniques to handle the merged electron-positron pairs, particularly in the case of ALP production via Higgs boson decays [8,15].

In this Letter, a search is presented for ALPs with a mass in the range 10 to 100 MeV and a proper decay length ($c\tau$) below 100 μm. The ALPs are produced in Higgs boson decays, and then, each ALP decays into an electron-positron pair ($H \to AA \to 4e$).

The analysis strategy is to, first, reconstruct the ALP candidates from merged electron-positron pairs. Next, we reconstruct a Higgs boson candidate from a pair of ALP candidates and search for a localized excess in the

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[*] Full author list given at the end of the Letter.

four-electron invariant mass spectrum near the known Higgs boson mass. The signal of a Higgs boson decaying to ALPs would manifest as a narrow resonant peak in the four-electron invariant mass distribution on top of a smoothly falling and predominantly nonresonant background.

The study is based on data from proton-proton (pp) collisions at a center-of-mass energy $\sqrt{s} = 13$ TeV. The data were collected with the CMS detector [19,20] at the LHC from 2016 to 2018, corresponding to an integrated luminosity of 138 fb^{-1}. Tabulated results are provided in the HEPData record for this analysis [21]. The full likelihood for this analysis is provided in [22].

The CMS apparatus [19,20] is a multipurpose, nearly hermetic detector, designed to trigger events of interest using a two-tiered trigger system [23,24] and identify electrons, muons, photons, and charged and neutral hadrons [25–27]. A global particle-flow (PF) algorithm [28] aims to reconstruct all individual particles in an event (PF candidates) by combining information from multiple subdetectors: the all-silicon inner tracker, the crystal electromagnetic calorimeter (ECAL), and the brass and scintillator hadron calorimeter (HCAL), all operating inside a 3.8 T superconducting solenoid, together with the gas-ionization muon system embedded in the flux-return yoke outside the solenoid. The PF candidates are then used to build τ leptons, jets, and missing transverse momentum [29–31]. Forward calorimeters, made of steel and quartz fibers, extend the pseudorapidity coverage provided by the barrel and end cap detectors. More detailed descriptions of the CMS detector, as well as the coordinate system and relevant kinematic variables, can be found in Refs. [19,20].

In the Monte Carlo (MC) simulation, the signal is generated using Higgs boson production via the gluon-fusion process, modeled with the POWHEG v2.0 generator at next-to-leading order (NLO) accuracy in quantum chromodynamics for the hard-scattering matrix element [32–35]. Parton showering, hadronization, and the Higgs boson decay are simulated with PYTHIA 8.240 [36], including underlying event modeling with the CP5 tune [37]. The parton distribution function set is NNPDF3.1 at next-to-next-to-leading order (NNLO) accuracy [38]. Higgs boson-related backgrounds include gluon fusion production with decays into a pair of photons (FxFx prescription [39] for jet matching and merging) as well as the electron-positron-photon three-body Dalitz decay, modeled at NLO by MadGraph5_aMC@NLO v2.6.5 [40]. The SM Z boson production is modeled at NNLO by the POWHEG v2.0 generator with the MiNNLO$_{\mathrm{PS}}$ prescription [41,42]. Double-photon production is modeled at NLO with Sherpa 2.2.4 [43] using the NNPDF3.0 [44] PDF, while single photon plus jet ($\gamma +$ jets) is modeled with PYTHIA 8.240.

Events are required to contain at least two reconstructed merged electron pairs (MEPs). Accordingly, high-level trigger paths that select two objects with significant electromagnetic energy [45] are used.

Despite the application of a novel end-to-end machine learning based reconstruction of the merged diphoton signature [46], the diphoton search [17] was limited by the angular resolution of the CMS ECAL. In this search, we target ALPs decaying to electron-positron pairs. Unlike the diphoton channel, the two charged particles leave hits in the silicon tracker. The tracker has an angular resolution that is superior to the ECAL, which enables us to resolve highly collimated pairs and reach sensitivity to ALP masses as low as 10 MeV. However, the standard CMS reconstruction does not include a dedicated algorithm for highly collimated electron-positron pairs and will typically reconstruct them as a single electron or misidentify them as converted photons. To address this, we have developed a set of dedicated reconstruction and identification techniques for merged electron-positron pairs. Henceforth, we refer to these as MEPs, since the electrons and positrons involved are not experimentally distinguished in this analysis. The MEPs are first reconstructed by matching two Gaussian-sum filter (GSF) [25,47] tracks from the silicon tracker to the same ECAL cluster; each track, extrapolated to the ECAL surface, must satisfy $\Delta R(\text{track, ECAL cluster}) < 0.015$, where ΔR is the Euclidean distance in the η-ϕ plane. An ECAL cluster is a collection of energy deposits in the ECAL that contains the calorimeter shower from an electron or a photon [25]. The GSF tracks are required to have a transverse momentum $p_{\mathrm{T}} > 5$ GeV and fewer than two missing hits in the inner tracker. The associated ECAL cluster is required to have an energy larger than 25 GeV, which is above the thresholds for the triggers used in this analysis. The p_{T} of the reconstructed MEPs, obtained from the vector sum of the two GSF tracks, is further required to be consistent with the ECAL cluster p_{T} within 70%. An adaptive vertex fitter [48] is used to fit the vertex of the MEP using the two GSF tracks. The resulting χ^2 value of the vertex fit is required to be smaller than five to reject failures of the vertex reconstruction. Successfully reconstructed MEPs are required to pass an isolation requirement, defined as the ratio of the summed transverse momentum and energy from the surrounding additional PF candidates within $0.015 < \Delta R < 0.3$. The isolation variable combines tracker, ECAL, and HCAL components, with the cutoff set at the 95% signal efficiency, to reject jets misidentified as MEPs.

A gradient boosted decision tree (GBDT) classifier, implemented with the XGBoost package [49], has been developed to identify MEPs that pass the isolation criteria and to discriminate ALP-induced MEPs from background objects. The backgrounds mainly come from photon conversions, electrons with additional tracks falling within the $\Delta R < 0.015$ cone centered on the electron momentum, and processes with jets that mimic the MEP signature. For example, processes such as Z boson, diphoton, and $\gamma +$ jets production can contribute as backgrounds. The input variables include observables related to the MEPs, such as the shower shape variables from the ECAL, and tracking-related variables such as the angular separation between the two matched GSF tracks, as well as the displacement of the reconstructed vertex with respect to the primary vertex (PV). The PV is taken to be the vertex corresponding to the hardest scattering in the event, evaluated using tracking information alone, as described in Section IX.4.1 of Ref. [50]. We exclude input variables that are explicitly correlated with the energy scale of the ALP candidate, such as p_T and energy, to avoid biasing ALP identification to a specific kinematic region of the MEP. The model is then trained using samples with different ALP mass hypotheses. The identification model is trained separately for four different data-taking periods in 2016 (split into two subperiods because of a change in the detector conditions), 2017, and 2018.

The identification model is trained using simulated samples of ALP signals, with various mass hypotheses, and simulated background samples, which include misreconstructed MEP candidates originating from photon conversions and misidentified electrons. The available samples are split into statistically independent subsets, with 80% used for training and 20% reserved for testing.

The GBDT working point (WP) chosen in the analysis, with a 70% signal efficiency for the classifier, reduces about 75% of MEP candidates from $Z \to ee$, about 55% from $\gamma +$ jets, and about 60% from diphoton events, where the efficiencies are defined as the ratio between the number of MEPs passing the GBDT WP and the number of MEPs passing the isolation requirement. This WP optimizes the expected sensitivity reach and yields a sufficient number of events for the statistical analysis. The performance of this identification WP is further corrected with a $Z \to \mu\mu\gamma$ control sample, where a pure sample of converted photons is used as a proxy for MEPs. Data from different data-taking eras are separated to validate the MEP selection and extract scale factors (SFs), defined as corrections for efficiency differences in MEP identification between data and simulation. This is accomplished by fitting the Z boson peak in the three-body invariant mass distribution of the $\mu\mu\gamma$ system in both data and simulation.

The $Z \to \mu\mu\gamma$ events used for this procedure are selected using standard CMS muon identification criteria [26] together with the Z boson mass constraint, ensuring a high-purity sample of photon conversions reconstructed as MEP candidates [25]. The same selection is applied to simulated events to extract the efficiency correction factors.

The SFs are consistent with unity and are checked by varying the signal and background functions for the fit or the mass range of the fit. They are also checked in different bins of vertex transverse displacement with respect to the PV and the ΔR distance between the two GSF tracks, as motivated by the kinematic differences between the signal and background processes. The systematic shifts from the former are under 10% for all cases, while for the latter, a closure within 5% is found when measuring the SFs in different regions of the ΔR between the two GSF tracks. The efficiencies of the trigger paths are evaluated using $Z \to ee$ events that pass the MEP identification criteria used for the event selection. The single-leg efficiency is found to be 95%, resulting in an overall event-level efficiency of about 90%.

The Higgs boson candidates are reconstructed from the two successfully identified MEPs, with the reconstructed candidate mass m_{4e} found from the invariant mass of the system formed by the two MEPs, assumed to be massless. This approximation is justified since the ALP mass is negligible compared to the Higgs boson mass and detector resolution.

The MEP energy values are determined by the merged ECAL clusters. The observed yield of the signal, for the process $H \to AA \to 4e$, is then extracted from this mass spectrum by fitting the data points to a set of parametric shapes, for which different modeling methods are used for resonant and nonresonant processes. For example, the expected number of signal events after the full selection varies from five to ten, depending on the ALP mass and lifetime hypothesis, assuming a branching fraction of around 5×10^{-5} corresponding to the observed sensitivity reach.

For the continuous background processes that originate from either prompt or nonprompt photon production and do not have a peaking structure around the Higgs boson mass, the falling shape of the mass spectrum is modeled from data. The background shape is determined by fitting the data in the sidebands of the 125 GeV Higgs boson mass peak, which are set as from 100 to 180 GeV with the 115 to 135 GeV window of the Higgs boson mass excluded. A set of F tests [51] is used to constrain the model complexity and determine the nominal parameters for the final fit. The candidate background functions include polynomial functions, exponential functions, Bernstein polynomials, Laurent series, and power-law functions, which are widely used in the CMS $H \to \gamma\gamma$ analyses, such as in Refs. [52,53]. The discrete profiling method [54] is introduced to account for the uncertainty of the parametric shape choices.

In addition to the nonresonant processes described above, which dominate the background, there is a nonnegligible resonant background. This is primarily from Higgs bosons decaying into two photons and the photons converting into electron-positron pairs, similar to ALPs. These converted photons are also reconstructed as MEPs and, therefore, produce a resonant background structure near the Higgs boson mass. This resonant background and the signal are modeled in approximately the same way. For both of them, a set of parametric shapes is used to model the peak structures from the intermediate Higgs boson resonance. The parametric shapes are fitted to simulated events that fulfill the event selection criteria. The parametric shape chosen for the peaking background is either a sum of multiple Gaussian functions, or a double-sided Crystal Ball function [55,56], with the choice depending on the number of events, in order to ensure fit stability. While the former can be used with a sufficient number of events to exploit a more complicated signal shape behavior, the latter is used with a limited event count to ensure a single peak structure.

The dominant systematic uncertainties come from the SFs, affecting the identification and trigger efficiencies. The uncertainties from the SF fitting procedure and the phase space closure are already accounted for, as discussed above. The nonresonant background shape uncertainty is handled via the discrete profiling method discussed above. The systematic uncertainties on the shape of the distributions for the signal and resonant background processes arise from the ECAL cluster energy scale and resolution corrections, evaluated using varied simulation templates. The uncertainty on the integrated luminosity applies to the yield of simulated processes, ranging from 1.2% to 2.5% depending on the data-taking year, while the overall 2016–2018 uncertainty is 1.6% [57–59]. Finally, the theoretical uncertainty in the Higgs boson production cross section [60] is included.

The statistical analysis is performed with an unbinned profile-likelihood fit to the four-electron invariant mass spectrum using asymptotic formulas, implemented with the Combine v8.2.0 package [61]. The measured distribution of the four-electron invariant mass is compared to the signal-plus-background fit in Fig. 1. No significant excess is observed over the background prediction. The observed limits are generally slightly higher than the expected limits, consistent with a small upward fluctuation of the data relative to the background-only prediction. The deviation remains within the 1-standard-deviation uncertainty band and, therefore, is statistically compatible with the background-only hypothesis.

Upper limits at 95% confidence level (CL) on the branching fraction for $H \to AA \to 4e$ are reported in Fig. 2. Figure 2 shows the 95% CL upper branching fraction limits as functions of the ALP mass for $c\tau = 1$, 10, and 100 μm, as well as the map of these limits in the ALP mass versus lifetime plane. For the last plot, points with the additional lifetime hypotheses are created by reweighting the existing samples, using the exponential probability distribution of the decay probability. The upper limit on the decay branching fraction is observed to be as low as $\mathcal{O}(10^{-5})$ at 95% CL, in the mass range of 10 to 100 MeV and for proper decay lengths between 0 and 10 μm. The sensitivity degrades at larger masses because of wider opening angles of the electron pairs and, also, at longer lifetimes because more ALPs start to fail the number of inner tracker hits requirement. An interpretation of these results in the effective ALP coupling model can be found in the Appendix of the End Matter.

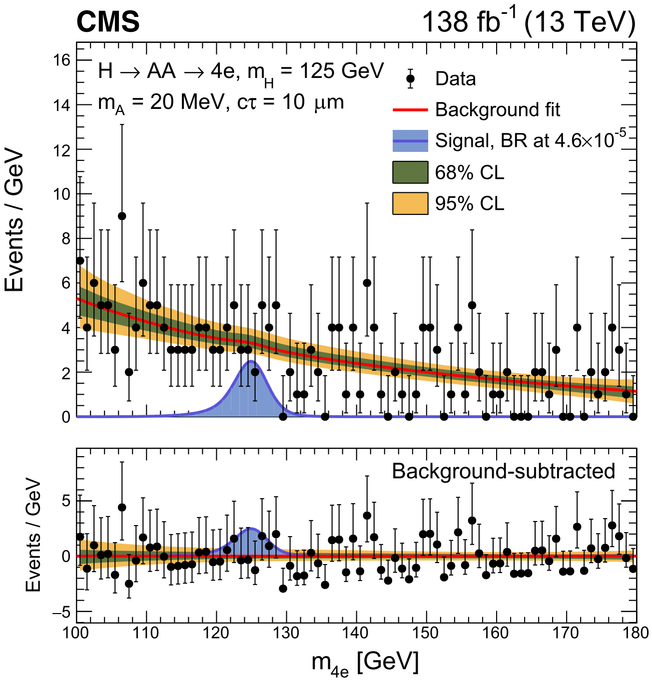

FIG. 1. Invariant mass distribution of the four-electron system (m_{4e}) for selected events (points), compared to the background-only fit (red) with its 68% and 95% CL uncertainty bands (green and yellow). A nonstacked benchmark signal (blue) for a Higgs boson decaying to a pair of ALPs with $m_A = 20$ MeV and $c\tau = 10$ μm is overlaid and normalized to a branching ratio (BR) of 4.6×10^{-5}, which corresponds to the 95% CL upper limit value set by this analysis. The lower panel shows the same data after subtracting the background fit.

This analysis establishes the first direct limits on the Higgs boson exotic decay $H \to AA \to 4e$ for an axionlike particle A with a mass of $\mathcal{O}(10)$ MeV, reaching branching fraction sensitivities as low as 10^{-5}. The limiting factor for the sensitivity of this analysis is the number of events in data, while the leading systematic uncertainty arises from the identification efficiency of merged electron-positron pairs. This search explores previously inaccessible parameter space and provides the most stringent constraints to date on this model. This search significantly improves the experimental sensitivity to axionlike particles with masses below 100 MeV, extending collider coverage to masses as low as 10 MeV for the first time. These results establish a new benchmark for future searches at the LHC.

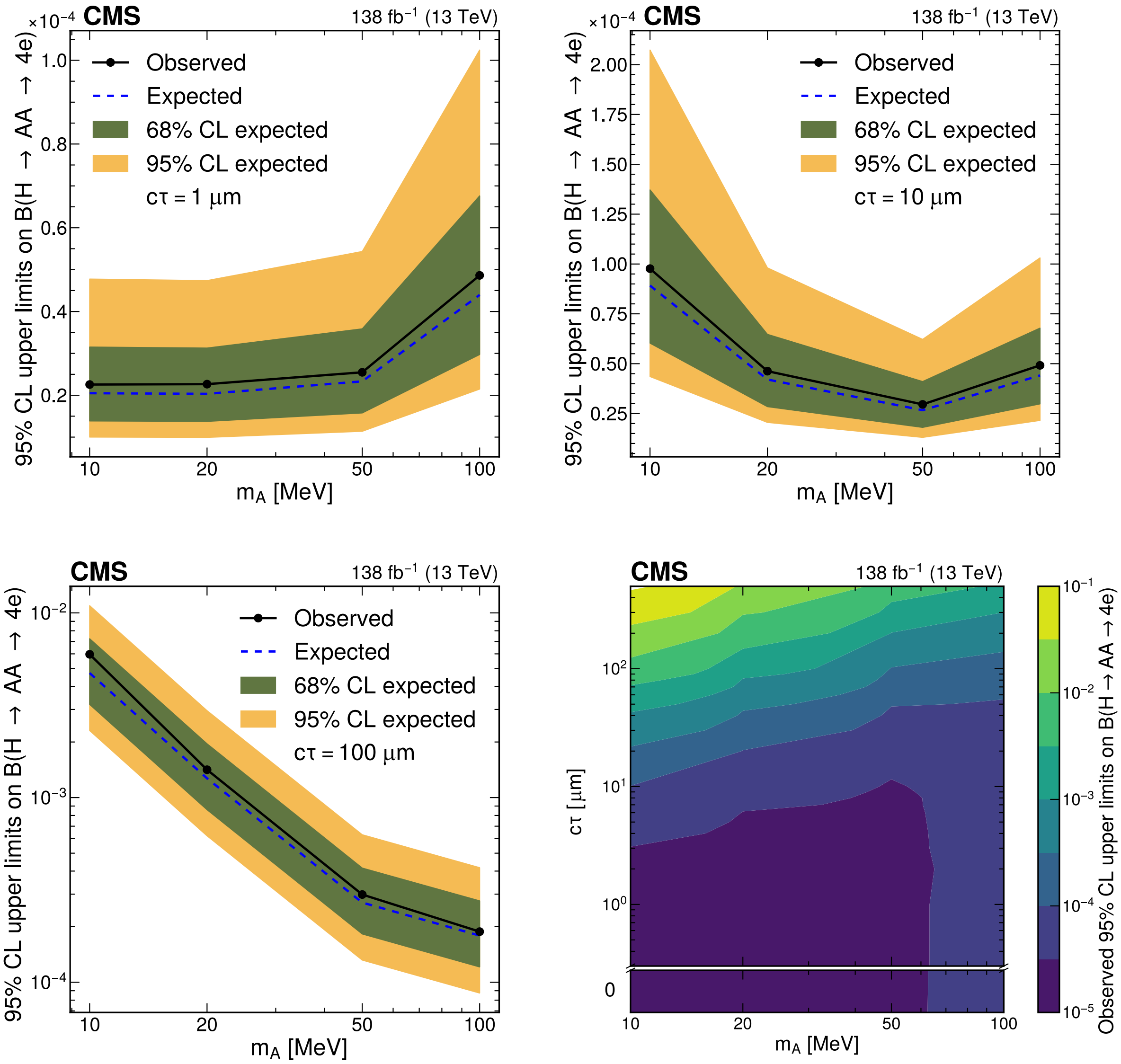

FIG. 2. Observed (solid points) and expected (dashed lines) 95% CL upper limits on the Higgs boson branching fraction to a pair of ALPs decaying into electron-positron pairs ($H \to AA \to ee$), shown as a function of the ALP mass for benchmark proper decay lengths of 1 μm (upper left), 10 μm (upper right), and 100 μm (lower left). The green and yellow bands represent the one and two standard deviation confidence intervals around the expected limits. The lower right panel shows a map of the observed 95% CL upper limit, shown as a color scale, as a function of the ALP mass m_A, and proper decay length $c\tau$.

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End Matter

Appendix: Interpretation of results in the effective ALP-Higgs coupling model—To bridge the gap between the collider and intensity-frontier experiments, the reported limits are interpreted in an effective ALP coupling model, as described in Refs. [15,63]. In this framework, the ALP couples to SM fermions and the Higgs doublet through derivative interactions. The effective Lagrangian terms governing these interactions are described in Ref. [15], and the corresponding partial decay widths for the ALP decaying into fermion pairs and for the Higgs boson decaying into a pair of ALPs are, at LO, given by

$$\Gamma(A \to f\bar{f}) = \frac{m_A m_f^2}{8\pi\Lambda^2} c_{ff}^2 \sqrt{1 - \frac{4m_f^2}{m_A^2}},$$
$$\Gamma(H \to AA) = \frac{v^2 m_H^3}{32\pi\Lambda^4} C_{AH}^2 \left(1 - \frac{2m_A^2}{m_H^2}\right)^2 \sqrt{1 - \frac{4m_A^2}{m_H^2}}. \tag{A1}$$

Here, c_{ff} is the strength of the flavor-dependent coupling between the ALP and an SM fermion, C_{AH} is the strength of the coupling between the ALP and the Higgs field, Λ is the characteristic energy scale of such an effective theory (often denoted as f_A in the literature on ALPs), v is the vacuum expectation value of the Higgs field, and m_A, m_f, and m_H denote the masses of the ALP, fermion, and Higgs boson, respectively.

With the assumption that ALPs only decay to electron-positron pairs, the total decay width of the ALP is entirely determined by the partial width in this channel, thus, relating its lifetime to the parameter c_{ff}/Λ. Therefore, the results reported in Fig. 2 can be reinterpreted as a function of this parameter, as shown in Fig. 3.

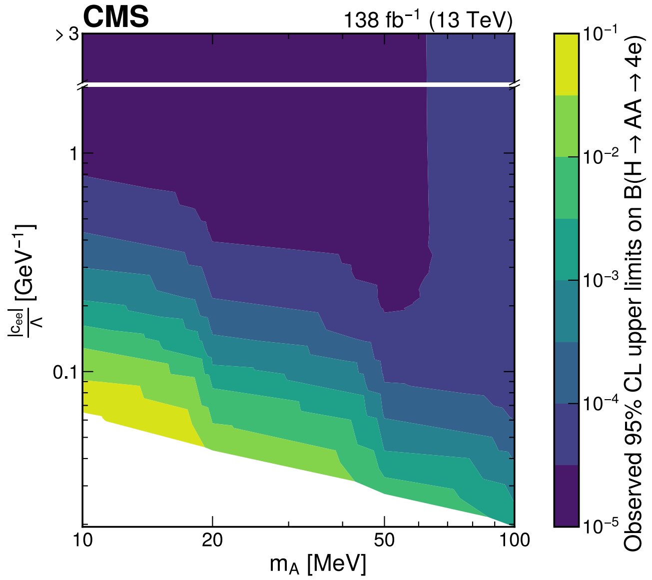

FIG. 3. A map of the observed 95% CL upper limit on the Higgs boson branching fraction for $H \to AA \to 4e$, as a function of the ALP mass and the ratio of the ALP coupling to electrons to the energy scale of the ALP effective interaction.

* * *

A. Hayrapetyan,[1] V. Makarenko,[1] A. Tumasyan,[1,b] W. Adam,[2] J. W. Andrejkovic,[2] L. Benato,[2] T. Bergauer,[2]
M. Dragicevic,[2] C. Giordano,[2] P. S. Hussain,[2] M. Jeitler,[2,c] N. Krammer,[2] A. Li,[2] D. Liko,[2] M. Matthewman,[2]
I. Mikulec,[2] J. Schieck,[2,c] R. Schöfbeck,[2,c] D. Schwarz,[2] M. Shooshtari,[2] M. Sonawane,[2] W. Waltenberger,[2]
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G. Bauer,[22,l,m] Z. Cui,[22,m] B. Li,[22,n] H. Wang,[22] K. Yi,[22,o] J. Zhang,[22] Y. Li,[23] Z. Lin,[24] C. Lu,[24] M. Xiao,[24,p]
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I. Salazar Segovia,[156] J. S. Tafoya Vargas,[156] W. Wei,[156] S. Yoo,[156] K. Adamidis,[157] M. Bachtis,[157] D. Campos,[157]
R. Cousins,[157] A. Datta,[157] G. Flores Avila,[157] J. Hauser,[157] M. Ignatenko,[157] M. A. Iqbal,[157] T. Lam,[157]
Y. f. Lo,[157] E. Manca,[157] A. Nunez Del Prado,[157] D. Saltzberg,[157] V. Valuev,[157] R. Clare,[158] J. W. Gary,[158]
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D. Diaz,[159] J. Duarte,[159] L. Giannini,[159] Y. Gu,[159] J. Guiang,[159] V. Krutelyov,[159] R. Lee,[159] J. Letts,[159] H. Li,[159]
M. Masciovecchio,[159] F. Mokhtar,[159] S. Mukherjee,[159] M. Pieri,[159] D. Primosch,[159] M. Quinnan,[159] V. Sharma,[159]
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C. Campagnari,[160] S. Carron Montero,[160,cccc] K. Downham,[160] C. Grieco,[160] M. M. Hussain,[160] J. Incandela,[160]
M. W. K. Lai,[160] A. J. Li,[160] P. Masterson,[160] J. Richman,[160] S. N. Santpur,[160] U. Sarica,[160] R. Schmitz,[160]
F. Setti,[160] J. Sheplock,[160] D. Stuart,[160] T. Á. Vámi,[160] X. Yan,[160] D. Zhang,[160] A. Albert,[161]
S. Bhattacharya,[161] A. Bornheim,[161] O. Cerri,[161] R. Kansal,[161] J. Mao,[161] H. B. Newman,[161] G. Reales Gutiérrez,[161]
T. Sievert,[161] M. Spiropulu,[161] J. R. Vlimant,[161] R. A. Wynne,[161] S. Xie,[161] J. Alison,[162] S. An,[162]
M. Cremonesi,[162] V. Dutta,[162] E. Y. Ertorer,[162] T. Ferguson,[162] T. A. Gómez Espinosa,[162] A. Harilal,[162]
A. Kallil Tharayil,[162] M. Kanemura,[162] C. Liu,[162] M. Marchegiani,[162] P. Meiring,[162] T. Mudholkar,[162] S. Murthy,[162]
P. Palit,[162] K. Park,[162] M. Paulini,[162] A. Roberts,[162] A. Sanchez,[162] W. Terrill,[162] J. P. Cumalat,[163]
W. T. Ford,[163] A. Hart,[163] S. Kwan,[163] J. Pearkes,[163] C. Savard,[163] N. Schonbeck,[163] K. Stenson,[163]
K. A. Ulmer,[163] S. R. Wagner,[163] N. Zipper,[163] D. Zuolo,[163] J. Alexander,[164] X. Chen,[164] J. Dickinson,[164]
A. Duquette,[164] J. Fan,[164] X. Fan,[164] J. Grassi,[164] S. Hogan,[164] P. Kotamnives,[164] J. Monroy,[164] G. Niendorf,[164]

M. Oshiro,[164] J. R. Patterson,[164] A. Ryd,[164] J. Thom,[164] P. Wittich,[164] R. Zou,[164] L. Zygala,[164] M. Albrow,[165] M. Alyari,[165] O. Amram,[165] G. Apollinari,[165] A. Apresyan,[165] L. A. T. Bauerdick,[165] D. Berry,[165] J. Berryhill,[165] P. C. Bhat,[165] K. Burkett,[165] J. N. Butler,[165] A. Canepa,[165] G. B. Cerati,[165] H. W. K. Cheung,[165] F. Chlebana,[165] C. Cosby,[165] G. Cummings,[165] I. Dutta,[165] V. D. Elvira,[165] J. Freeman,[165] A. Gandrakota,[165] Z. Gecse,[165] L. Gray,[165] D. Green,[165] A. Grummer,[165] S. Grünendahl,[165] D. Guerrero,[165] O. Gutsche,[165] R. M. Harris,[165] T. C. Herwig,[165] J. Hirschauer,[165] V. Innocente,[165] B. Jayatilaka,[165] S. Jindariani,[165] M. Johnson,[165] U. Joshi,[165] B. Klima,[165] K. H. M. Kwok,[165] S. Lammel,[165] C. Lee,[165] D. Lincoln,[165] R. Lipton,[165] T. Liu,[165] K. Maeshima,[165] D. Mason,[165] P. McBride,[165] P. Merkel,[165] S. Mrenna,[165] S. Nahn,[165] J. Ngadiuba,[165] D. Noonan,[165] S. Norberg,[165] V. Papadimitriou,[165] N. Pastika,[165] K. Pedro,[165] C. Pena,[165,dddd] C. E. Perez Lara,[165] F. Ravera,[165] A. Reinsvold Hall,[165,eeee] L. Ristori,[165] M. Safdari,[165] E. Sexton-Kennedy,[165] N. Smith,[165] A. Soha,[165] L. Spiegel,[165] S. Stoynev,[165] J. Strait,[165] L. Taylor,[165] S. Tkaczyk,[165] N. V. Tran,[165] L. Uplegger,[165] E. W. Vaandering,[165] C. Wang,[165] I. Zoi,[165] C. Aruta,[166] P. Avery,[166] D. Bourilkov,[166] P. Chang,[166] V. Cherepanov,[166] R. D. Field,[166] C. Huh,[166] E. Koenig,[166] M. Kolosova,[166] J. Konigsberg,[166] A. Korytov,[166] G. Mitselmakher,[166] K. Mohrman,[166] A. Muthirakalayil Madhu,[166] N. Rawal,[166] S. Rosenzweig,[166] V. Sulimov,[166] Y. Takahashi,[166] J. Wang,[166] T. Adams,[167] A. Al Kadhim,[167] A. Askew,[167] S. Bower,[167] R. Goff,[167] R. Hashmi,[167] A. Hassani,[167] R. S. Kim,[167] T. Kolberg,[167] G. Martinez,[167] M. Mazza,[167] H. Prosper,[167] P. R. Prova,[167] R. Yohay,[167] B. Alsufyani,[168] S. Butalla,[168] S. Das,[168] M. Hohlmann,[168] M. Lavinsky,[168] E. Yanes,[168] M. R. Adams,[169] N. Barnett,[169] A. Baty,[169] C. Bennett,[169] R. Cavanaugh,[169] R. Escobar Franco,[169] O. Evdokimov,[169] C. E. Gerber,[169] H. Gupta,[169] M. Hawksworth,[169] A. Hingrajiya,[169] D. J. Hofman,[169] Z. Huang,[169] J. h. Lee,[169] C. Mills,[169] S. Nanda,[169] G. Nigmatkulov,[169] B. Ozek,[169] T. Phan,[169] D. Pilipovic,[169] R. Pradhan,[169] E. Prifti,[169] P. Roy,[169] T. Roy,[169] N. Singh,[169] M. B. Tonjes,[169] N. Varelas,[169] M. A. Wadud,[169] J. Yoo,[169] M. Alhusseini,[170] D. Blend,[170] K. Dilsiz,[170,ffff] O. K. Köseyan,[170] A. Mestvirishvili,[170,gggg] O. Neogi,[170] H. Ogul,[170,hhhh] Y. Onel,[170] A. Penzo,[170] C. Snyder,[170] E. Tiras,[170,iiii] B. Blumenfeld,[171] J. Davis,[171] A. V. Gritsan,[171] L. Kang,[171] S. Kyriacou,[171] P. Maksimovic,[171] M. Roguljic,[171] S. Sekhar,[171] M. V. Srivastav,[171] M. Swartz,[171] A. Abreu,[172] L. F. Alcerro Alcerro,[172] J. Anguiano,[172] S. Arteaga Escatel,[172] P. Baringer,[172] A. Bean,[172] R. Bhattacharya,[172] Z. Flowers,[172] D. Grove,[172] J. King,[172] G. Krintiras,[172] M. Lazarovits,[172] C. Le Mahieu,[172] J. Marquez,[172] M. Murray,[172] M. Nickel,[172] S. Popescu,[172,jjjj] C. Rogan,[172] C. Royon,[172] S. Rudrabhatla,[172] S. Sanders,[172] C. Smith,[172] G. Wilson,[172] B. Allmond,[173] N. Islam,[173] A. Ivanov,[173] K. Kaadze,[173] Y. Maravin,[173] J. Natoli,[173] G. G. Reddy,[173] D. Roy,[173] G. Sorrentino,[173] A. Baden,[174] A. Belloni,[174] J. Bistany-riebman,[174] S. C. Eno,[174] N. J. Hadley,[174] S. Jabeen,[174] R. G. Kellogg,[174] T. Koeth,[174] B. Kronheim,[174] S. Lascio,[174] P. Major,[174] A. C. Mignerey,[174] C. Palmer,[174] C. Papageorgakis,[174] M. M. Paranjpe,[174] E. Popova,[174,kkkk] A. Shevelev,[174] L. Zhang,[174] C. Baldenegro Barrera,[175] H. Bossi,[175] S. Bright-Thonney,[175] I. A. Cali,[175] Y. c. Chen,[175] P. c. Chou,[175] M. D'Alfonso,[175] J. Eysermans,[175] C. Freer,[175] G. Gomez-Ceballos,[175] M. Goncharov,[175] G. Grosso,[175] P. Harris,[175] D. Hoang,[175] G. M. Innocenti,[175] K. Ivanov,[175] D. Kovalskyi,[175] J. Krupa,[175] L. Lavezzo,[175] Y.-J. Lee,[175] K. Long,[175] C. Mcginn,[175] A. Novak,[175] M. I. Park,[175] C. Paus,[175] C. Reissel,[175] C. Roland,[175] G. Roland,[175] S. Rothman,[175] T. a. Sheng,[175] G. S. F. Stephans,[175] D. Walter,[175] J. Wang,[175] Z. Wang,[175] B. Wyslouch,[175] T. J. Yang,[175] B. Crossman,[176] W. J. Jackson,[176] C. Kapsiak,[176] M. Krohn,[176] D. Mahon,[176] J. Mans,[176] B. Marzocchi,[176] R. Rusack,[176] O. Sancar,[176] R. Saradhy,[176] N. Strobbe,[176] K. Bloom,[177] D. R. Claes,[177] G. Haza,[177] J. Hossain,[177] C. Joo,[177] I. Kravchenko,[177] A. Rohilla,[177] J. E. Siado,[177] W. Tabb,[177] A. Vagnerini,[177] A. Wightman,[177] F. Yan,[177] H. Bandyopadhyay,[178] L. Hay,[178] H. w. Hsia,[178] I. Iashvili,[178] A. Kalogeropoulos,[178] A. Kharchilava,[178] A. Mandal,[178] M. Morris,[178] D. Nguyen,[178] S. Rappoccio,[178] H. Rejeb Sfar,[178] A. Williams,[178] P. Young,[178] D. Yu,[178] G. Alverson,[179] E. Barberis,[179] J. Bonilla,[179] B. Bylsma,[179] M. Campana,[179] J. Dervan,[179] Y. Haddad,[179] Y. Han,[179] I. Israr,[179] A. Krishna,[179] M. Lu,[179] N. Manganelli,[179] R. Mccarthy,[179] D. M. Morse,[179] T. Orimoto,[179] L. Skinnari,[179] C. S. Thoreson,[179] E. Tsai,[179] D. Wood,[179] S. Dittmer,[180] K. A. Hahn,[180] M. Mcginnis,[180] Y. Miao,[180] D. G. Monk,[180] M. H. Schmitt,[180] A. Taliercio,[180] M. Velasco,[180] J. Wang,[180] G. Agarwal,[181] R. Band,[181] R. Bucci,[181] S. Castells,[181] A. Das,[181] A. Ehnis,[181] R. Goldouzian,[181] M. Hildreth,[181] K. Hurtado Anampa,[181] T. Ivanov,[181] C. Jessop,[181] A. Karneyeu,[181] K. Lannon,[181] J. Lawrence,[181] N. Loukas,[181] L. Lutton,[181] J. Mariano,[181] N. Marinelli,[181]

I. Mcalister,[181] T. McCauley,[181] C. Mcgrady,[181] C. Moore,[181] Y. Musienko,[181,kkkk] H. Nelson,[181] M. Osherson,[181]
A. Piccinelli,[181] R. Ruchti,[181] A. Townsend,[181] Y. Wan,[181] M. Wayne,[181] H. Yockey,[181] A. Basnet,[182]
M. Carrigan,[182] R. De Los Santos,[182] L. S. Durkin,[182] C. Hill,[182] M. Joyce,[182] M. Nunez Ornelas,[182]
D. A. Wenzl,[182] B. L. Winer,[182] B. R. Yates,[182] H. Bouchamaoui,[183] G. Dezoort,[183] P. Elmer,[183] A. Frankenthal,[183]
M. Galli,[183] B. Greenberg,[183] N. Haubrich,[183] K. Kennedy,[183] G. Kopp,[183] Y. Lai,[183] D. Lange,[183]
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D. Kondratyev,[185] M. Liu,[185] G. Negro,[185] N. Neumeister,[185] G. Paspalaki,[185] S. Piperov,[185] N. R. Saha,[185]
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E. Shumka,[186] D. Acosta,[187] A. Agrawal,[187] C. Arbour,[187] T. Carnahan,[187] P. Das,[187] K. M. Ecklund,[187]
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O. Miguel Colin,[187] B. P. Padley,[187] R. Redjimi,[187] J. Rotter,[187] M. Wulansatiti,[187] E. Yigitbasi,[187] Y. Zhang,[187]
O. Bessidskaia Bylund,[188] A. Bodek,[188] P. de Barbaro,[188,a] R. Demina,[188] A. Garcia-Bellido,[188] H. S. Hare,[188]
O. Hindrichs,[188] N. Parmar,[188] P. Parygin,[188,kkkk] H. Seo,[188] R. Taus,[188] B. Chiarito,[189] J. P. Chou,[189]
S. V. Clark,[189] S. Donnelly,[189] D. Gadkari,[189] Y. Gershtein,[189] E. Halkiadakis,[189] C. Houghton,[189]
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S. Thomas,[189] J. Vora,[189] D. Ally,[190] A. G. Delannoy,[190] S. Fiorendi,[190] J. Harris,[190] T. Holmes,[190]
A. R. Kanuganti,[190] N. Karunarathna,[190] J. Lawless,[190] L. Lee,[190] E. Nibigira,[190] B. Skipworth,[190] S. Spanier,[190]
D. Aebi,[191] M. Ahmad,[191] T. Akhter,[191] K. Androsov,[191] A. Bolshov,[191] O. Bouhali,[191,llll] A. Cagnotta,[191]
V. D'Amante,[191] R. Eusebi,[191] P. Flanagan,[191] J. Gilmore,[191] Y. Guo,[191] T. Kamon,[191] S. Luo,[191] R. Mueller,[191]
A. Safonov,[191] N. Akchurin,[192] J. Damgov,[192] Y. Feng,[192] N. Gogate,[192] Y. Kazhykarim,[192] K. Lamichhane,[192]
S. W. Lee,[192] C. Madrid,[192] A. Mankel,[192] T. Peltola,[192] I. Volobouev,[192] E. Appelt,[193] Y. Chen,[193] S. Greene,[193]
A. Gurrola,[193] W. Johns,[193] R. Kunnawalkam Elayavalli,[193] A. Melo,[193] D. Rathjens,[193] F. Romeo,[193]
P. Sheldon,[193] S. Tuo,[193] J. Velkovska,[193] J. Viinikainen,[193] J. Zhang,[193] B. Cardwell,[194] H. Chung,[194] B. Cox,[194]
J. Hakala,[194] G. Hamilton Ilha Machado,[194] R. Hirosky,[194] M. Jose,[194] A. Ledovskoy,[194] C. Mantilla,[194] C. Neu,[194]
C. Ramón Álvarez,[194] Z. Wu,[194] S. Bhattacharya,[195] P. E. Karchin,[195] A. Aravind,[196] S. Banerjee,[196] K. Black,[196]
T. Bose,[196] E. Chavez,[196] S. Dasu,[196] P. Everaerts,[196] C. Galloni,[196] H. He,[196] M. Herndon,[196] A. Herve,[196]
C. K. Koraka,[196] S. Lomte,[196] R. Loveless,[196] A. Mallampalli,[196] A. Mohammadi,[196] S. Mondal,[196] T. Nelson,[196]
G. Parida,[196] L. Pétré,[196] D. Pinna,[196] A. Savin,[196] V. Shang,[196] V. Sharma,[196] W. H. Smith,[196] D. Teague,[196]
H. F. Tsoi,[196] W. Vetens,[196] A. Warden,[196] S. Afanasiev,[197] V. Alexakhin,[197] Yu. Andreev,[197] T. Aushev,[197]
D. Budkouski,[197] R. Chistov,[197] M. Danilov,[197] T. Dimova,[197] A. Ershov,[197] S. Gninenko,[197] I. Gorbunov,[197]
A. Gribushin,[197] A. Kamenev,[197] V. Karjavine,[197] M. Kirsanov,[197] V. Klyukhin,[197] O. Kodolova,[197,mmmm]
V. Korenkov,[197] I. Korsakov,[197] A. Kozyrev,[197] N. Krasnikov,[197] A. Lanev,[197] A. Malakhov,[197] V. Matveev,[197]
A. Nikitenko,[197,nnnn,mmmm] V. Palichik,[197] V. Perelygin,[197] S. Petrushanko,[197] S. Polikarpov,[197] O. Radchenko,[197]
M. Savina,[197] V. Shalaev,[197] S. Shmatov,[197] S. Shulha,[197] Y. Skovpen,[197] K. Slizhevskiy,[197] V. Smirnov,[197]
O. Teryaev,[197] I. Tlisova,[197] A. Toropin,[197] N. Voytishin,[197] A. Zarubin,[197] I. Zhizhin,[197] L. Dudko,[198]
V. Kim,[198,kkkk] V. Murzin,[198] V. Oreshkin,[198] D. Sosnov,[198] E. Boos,[198] V. Bunichev,[198] M. Dubinin,[198,dddd]
V. Savrin,[198] and A. Snigirev[198]

(CMS Collaboration)

[1]*Yerevan Physics Institute, Yerevan, Armenia*
[2]*Institut für Hochenergiephysik, Vienna, Austria*
[3]*Universiteit Antwerpen, Antwerpen, Belgium*
[4]*Vrije Universiteit Brussel, Brussel, Belgium*
[5]*Université Libre de Bruxelles, Bruxelles, Belgium*
[6]*Ghent University, Ghent, Belgium*
[7]*Université Catholique de Louvain, Louvain-la-Neuve, Belgium*
[8]*Centro Brasileiro de Pesquisas Fisicas, Rio de Janeiro, Brazil*

[9]*Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil*
[10]*Universidade Estadual Paulista, Universidade Federal do ABC, São Paulo, Brazil*
[11]*Institute for Nuclear Research and Nuclear Energy, Bulgarian Academy of Sciences, Sofia, Bulgaria*
[12]*University of Sofia, Sofia, Bulgaria*
[13]*Instituto De Alta Investigación, Universidad de Tarapacá, Casilla 7 D, Arica, Chile*
[14]*Universidad Tecnica Federico Santa Maria, Valparaiso, Chile*
[15]*Beihang University, Beijing, China*
[16]*Department of Physics, Tsinghua University, Beijing, China*
[17]*Institute of High Energy Physics, Beijing, China*
[18]*State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China*
[19]*State Key Laboratory of Nuclear Physics and Technology, Institute of Quantum Matter, South China Normal University, Guangzhou, China*
[20]*Sun Yat-Sen University, Guangzhou, China*
[21]*University of Science and Technology of China, Hefei, China*
[22]*Nanjing Normal University, Nanjing, China*
[23]*Institute of Modern Physics and Key Laboratory of Nuclear Physics and Ion-beam Application (MOE) - Fudan University, Shanghai, China*
[24]*Zhejiang University, Hangzhou, Zhejiang, China*
[25]*Universidad de Los Andes, Bogota, Colombia*
[26]*Universidad de Antioquia, Medellin, Colombia*
[27]*University of Split, Faculty of Electrical Engineering, Mechanical Engineering and Naval Architecture, Split, Croatia*
[28]*University of Split, Faculty of Science, Split, Croatia*
[29]*Institute Rudjer Boskovic, Zagreb, Croatia*
[30]*University of Cyprus, Nicosia, Cyprus*
[31]*Charles University, Prague, Czech Republic*
[32]*Escuela Politecnica Nacional, Quito, Ecuador*
[33]*Universidad San Francisco de Quito, Quito, Ecuador*
[34]*Academy of Scientific Research and Technology of the Arab Republic of Egypt, Egyptian Network of High Energy Physics, Cairo, Egypt*
[35]*Center for High Energy Physics (CHEP-FU), Fayoum University, El-Fayoum, Egypt*
[36]*National Institute of Chemical Physics and Biophysics, Tallinn, Estonia*
[37]*Department of Physics, University of Helsinki, Helsinki, Finland*
[38]*Helsinki Institute of Physics, Helsinki, Finland*
[39]*Lappeenranta-Lahti University of Technology, Lappeenranta, Finland*
[40]*IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France*
[41]*Laboratoire Leprince-Ringuet, CNRS/IN2P3, Ecole Polytechnique, Institut Polytechnique de Paris, Palaiseau, France*
[42]*Université de Strasbourg, CNRS, IPHC UMR 7178, Strasbourg, France*
[43]*Centre de Calcul de l'Institut National de Physique Nucleaire et de Physique des Particules, CNRS/IN2P3, Villeurbanne, France*
[44]*Institut de Physique des 2 Infinis de Lyon (IP2I), Villeurbanne, France*
[45]*Georgian Technical University, Tbilisi, Georgia*
[46]*RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany*
[47]*RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany*
[48]*RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany*
[49]*Deutsches Elektronen-Synchrotron, Hamburg, Germany*
[50]*University of Hamburg, Hamburg, Germany*
[51]*Karlsruher Institut fuer Technologie, Karlsruhe, Germany*
[52]*Institute of Nuclear and Particle Physics (INPP), NCSR Demokritos, Aghia Paraskevi, Greece*
[53]*National and Kapodistrian University of Athens, Athens, Greece*
[54]*National Technical University of Athens, Athens, Greece*
[55]*University of Ioánnina, Ioánnina, Greece*
[56]*HUN-REN Wigner Research Centre for Physics, Budapest, Hungary*
[57]*MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary*
[58]*Faculty of Informatics, University of Debrecen, Debrecen, Hungary*
[59]*HUN-REN ATOMKI - Institute of Nuclear Research, Debrecen, Hungary*
[60]*Karoly Robert Campus, MATE Institute of Technology, Gyongyos, Hungary*
[61]*Panjab University, Chandigarh, India*
[62]*University of Delhi, Delhi, India*
[63]*Indian Institute of Technology Mandi (IIT-Mandi), Himachal Pradesh, India*
[64]*University of Hyderabad, Hyderabad, India*
[65]*Indian Institute of Technology Kanpur, Kanpur, India*

[66]*Saha Institute of Nuclear Physics, HBNI, Kolkata, India*
[67]*Indian Institute of Technology Madras, Madras, India*
[68]*IISER Mohali, India, Mohali, India*
[69]*Tata Institute of Fundamental Research-A, Mumbai, India*
[70]*Tata Institute of Fundamental Research-B, Mumbai, India*
[71]*National Institute of Science Education and Research, An OCC of Homi Bhabha National Institute, Bhubaneswar, Odisha, India*
[72]*Indian Institute of Science Education and Research (IISER), Pune, India*
[73]*Indian Institute of Technology Hyderabad, Telangana, India*
[74]*Isfahan University of Technology, Isfahan, Iran*
[75]*Institute for Research in Fundamental Sciences (IPM), Tehran, Iran*
[76]*University College Dublin, Dublin, Ireland*
[77a]*INFN Sezione di Bari, Bari, Italy*
[77b]*Università di Bari, Bari, Italy*
[77c]*Politecnico di Bari, Bari, Italy*
[78a]*INFN Sezione di Bologna, Bologna, Italy*
[78b]*Università di Bologna, Bologna, Italy*
[79a]*INFN Sezione di Catania, Catania, Italy*
[79b]*Università di Catania, Catania, Italy*
[80a]*INFN Sezione di Firenze, Firenze, Italy*
[80b]*Università di Firenze, Firenze, Italy*
[81]*INFN Laboratori Nazionali di Frascati, Frascati, Italy*
[82a]*INFN Sezione di Genova, Genova, Italy*
[82b]*Università di Genova, Genova, Italy*
[83a]*INFN Sezione di Milano-Bicocca, Milano, Italy*
[83b]*Università di Milano-Bicocca, Milano, Italy*
[84a]*INFN Sezione di Napoli, Napoli, Italy*
[84b]*Università di Napoli 'Federico II', Napoli, Italy*
[84c]*Università della Basilicata, Potenza, Italy*
[84d]*Scuola Superiore Meridionale (SSM), Napoli, Italy*
[85a]*INFN Sezione di Padova, Padova, Italy*
[85b]*Università di Padova, Padova, Italy*
[85c]*Universita degli Studi di Cagliari, Cagliari, Italy*
[86a]*INFN Sezione di Pavia, Pavia, Italy*
[86b]*Università di Pavia, Pavia, Italy*
[87a]*INFN Sezione di Perugia, Perugia, Italy*
[87b]*Università di Perugia, Perugia, Italy*
[88a]*INFN Sezione di Pisa, Pisa, Italy*
[88b]*Università di Pisa, Pisa, Italy*
[88c]*Scuola Normale Superiore di Pisa, Pisa, Italy*
[88d]*Università di Siena, Siena, Italy*
[89a]*INFN Sezione di Roma, Roma, Italy*
[89b]*Sapienza Università di Roma, Roma, Italy*
[90a]*INFN Sezione di Torino, Torino, Italy*
[90b]*Università di Torino, Torino, Italy*
[90c]*Università del Piemonte Orientale, Novara, Italy*
[91a]*INFN Sezione di Trieste, Trieste, Italy*
[91b]*Università di Trieste, Trieste, Italy*
[92]*Kyungpook National University, Daegu, Korea*
[93]*Department of Mathematics and Physics - GWNU, Gangneung, Korea*
[94]*Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea*
[95]*Hanyang University, Seoul, Korea*
[96]*Korea University, Seoul, Korea*
[97]*Kyung Hee University, Department of Physics, Seoul, Korea*
[98]*Sejong University, Seoul, Korea*
[99]*Seoul National University, Seoul, Korea*
[100]*University of Seoul, Seoul, Korea*
[101]*Yonsei University, Department of Physics, Seoul, Korea*
[102]*Sungkyunkwan University, Suwon, Korea*
[103]*College of Engineering and Technology, American University of the Middle East (AUM), Dasman, Kuwait*
[104]*Kuwait University - College of Science - Department of Physics, Safat, Kuwait*

[105]*Riga Technical University, Riga, Latvia*
[106]*University of Latvia (LU), Riga, Latvia*
[107]*Vilnius University, Vilnius, Lithuania*
[108]*National Centre for Particle Physics, Universiti Malaya, Kuala Lumpur, Malaysia*
[109]*Universidad de Sonora (UNISON), Hermosillo, Mexico*
[110]*Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico*
[111]*Universidad Iberoamericana, Mexico City, Mexico*
[112]*Benemerita Universidad Autonoma de Puebla, Puebla, Mexico*
[113]*University of Montenegro, Podgorica, Montenegro*
[114]*University of Canterbury, Christchurch, New Zealand*
[115]*National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan*
[116]*AGH University of Krakow, Krakow, Poland*
[117]*National Centre for Nuclear Research, Swierk, Poland*
[118]*Institute of Experimental Physics, Faculty of Physics, University of Warsaw, Warsaw, Poland*
[119]*Warsaw University of Technology, Warsaw, Poland*
[120]*Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal*
[121]*Faculty of Physics, University of Belgrade, Belgrade, Serbia*
[122]*VINCA Institute of Nuclear Sciences, University of Belgrade, Belgrade, Serbia*
[123]*Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain*
[124]*Universidad Autónoma de Madrid, Madrid, Spain*
[125]*Universidad de Oviedo, Instituto Universitario de Ciencias y Tecnologías Espaciales de Asturias (ICTEA), Oviedo, Spain*
[126]*Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain*
[127]*University of Colombo, Colombo, Sri Lanka*
[128]*University of Ruhuna, Department of Physics, Matara, Sri Lanka*
[129]*CERN, European Organization for Nuclear Research, Geneva, Switzerland*
[130]*PSI Center for Neutron and Muon Sciences, Villigen, Switzerland*
[131]*ETH Zurich - Institute for Particle Physics and Astrophysics (IPA), Zurich, Switzerland*
[132]*Universität Zürich, Zurich, Switzerland*
[133]*National Central University, Chung-Li, Taiwan*
[134]*National Taiwan University (NTU), Taipei, Taiwan*
[135]*High Energy Physics Research Unit, Department of Physics, Faculty of Science, Chulalongkorn University, Bangkok, Thailand*
[136]*Tunis El Manar University, Tunis, Tunisia*
[137]*Çukurova University, Physics Department, Science and Art Faculty, Adana, Turkey*
[138]*Hacettepe University, Ankara, Turkey*
[139]*Middle East Technical University, Physics Department, Ankara, Turkey*
[140]*Bogazici University, Istanbul, Turkey*
[141]*Istanbul Technical University, Istanbul, Turkey*
[142]*Istanbul University, Istanbul, Turkey*
[143]*Yildiz Technical University, Istanbul, Turkey*
[144]*Institute for Scintillation Materials of National Academy of Science of Ukraine, Kharkiv, Ukraine*
[145]*National Science Centre, Kharkiv Institute of Physics and Technology, Kharkiv, Ukraine*
[146]*University of Bristol, Bristol, United Kingdom*
[147]*Rutherford Appleton Laboratory, Didcot, United Kingdom*
[148]*Imperial College, London, United Kingdom*
[149]*Brunel University, Uxbridge, United Kingdom*
[150]*Baylor University, Waco, Texas, USA*
[151]*Bethel University, St. Paul, Minnesota, USA*
[152]*Catholic University of America, Washington, DC, USA*
[153]*The University of Alabama, Tuscaloosa, Alabama, USA*
[154]*Boston University, Boston, Massachusetts, USA*
[155]*Brown University, Providence, Rhode Island, USA*
[156]*University of California, Davis, Davis, California, USA*
[157]*University of California, Los Angeles, California, USA*
[158]*University of California, Riverside, Riverside, California, USA*
[159]*University of California, San Diego, La Jolla, California, USA*
[160]*University of California, Santa Barbara - Department of Physics, Santa Barbara, California, USA*
[161]*California Institute of Technology, Pasadena, California, USA*
[162]*Carnegie Mellon University, Pittsburgh, Pennsylvania, USA*
[163]*University of Colorado Boulder, Boulder, Colorado, USA*
[164]*Cornell University, Ithaca, New York, USA*

[165]*Fermi National Accelerator Laboratory, Batavia, Illinois, USA*
[166]*University of Florida, Gainesville, Florida, USA*
[167]*Florida State University, Tallahassee, Florida, USA*
[168]*Florida Institute of Technology, Melbourne, Florida, USA*
[169]*University of Illinois Chicago, Chicago, Illinois, USA*
[170]*The University of Iowa, Iowa City, Iowa, USA*
[171]*Johns Hopkins University, Baltimore, Maryland, USA*
[172]*The University of Kansas, Lawrence, Kansas, USA*
[173]*Kansas State University, Manhattan, Kansas, USA*
[174]*University of Maryland, College Park, Maryland, USA*
[175]*Massachusetts Institute of Technology, Cambridge, Massachusetts, USA*
[176]*University of Minnesota, Minneapolis, Minnesota, USA*
[177]*University of Nebraska-Lincoln, Lincoln, Nebraska, USA*
[178]*State University of New York at Buffalo, Buffalo, New York, USA*
[179]*Northeastern University, Boston, Massachusetts, USA*
[180]*Northwestern University, Evanston, Illinois, USA*
[181]*University of Notre Dame, Notre Dame, Indiana, USA*
[182]*The Ohio State University, Columbus, Ohio, USA*
[183]*Princeton University, Princeton, New Jersey, USA*
[184]*University of Puerto Rico, Mayaguez, Puerto Rico, USA*
[185]*Purdue University, West Lafayette, Indiana, USA*
[186]*Purdue University Northwest, Hammond, Indiana, USA*
[187]*Rice University, Houston, Texas, USA*
[188]*University of Rochester, Rochester, New York, USA*
[189]*Rutgers, The State University of New Jersey, Piscataway, New Jersey, USA*
[190]*University of Tennessee, Knoxville, Tennessee, USA*
[191]*Texas A&M University, College Station, Texas, USA*
[192]*Texas Tech University, Lubbock, Texas, USA*
[193]*Vanderbilt University, Nashville, Tennessee, USA*
[194]*University of Virginia, Charlottesville, Virginia, USA*
[195]*Wayne State University, Detroit, Michigan, USA*
[196]*University of Wisconsin - Madison, Madison, Wisconsin, USA*
[197]*An institute or international laboratory covered by a cooperation agreement with CERN*
[198]*An institute formerly covered by a cooperation agreement with CERN*

[a] Deceased.
[b] Also at Yerevan State University, Yerevan, Armenia.
[c] Also at TU Wien, Vienna, Austria.
[d] Also at Ghent University, Ghent, Belgium.
[e] Also at FACAMP—Faculdades de Campinas, Sao Paulo, Brazil.
[f] Also at Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil.
[g] Also at Universidade Estadual de Campinas, Campinas, Brazil.
[h] Also at Federal University of Rio Grande do Sul, Porto Alegre, Brazil.
[i] Also at The University of the State of Amazonas, Manaus, Brazil.
[j] Also at University of Chinese Academy of Sciences, Beijing, China.
[k] Also at China Center of Advanced Science and Technology, Beijing, China.
[l] Also at School of Physics, Zhengzhou University, Zhengzhou, China.
[m] Also at Henan Normal University, Xinxiang, China.
[n] Also at University of Shanghai for Science and Technology, Shanghai, China.
[o] Also at The University of Iowa, Iowa City, Iowa, USA.
[p] Also at Center for High Energy Physics, Peking University, Beijing, China.
[q] Also at Helwan University, Cairo, Egypt.
[r] Also at Zewail City of Science and Technology, Zewail, Egypt.
[s] Also at British University in Egypt, Cairo, Egypt.
[t] Also at Purdue University, West Lafayette, Indiana, USA.
[u] Also at Université de Haute Alsace, Mulhouse, France.
[v] Also at Another institute or international laboratory covered by a cooperation agreement with CERN.
[w] Also at University of Hamburg, Hamburg, Germany.
[x] Also at RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany.
[y] Also at Bergische University Wuppertal (BUW), Wuppertal, Germany.

[z] Also at Brandenburg University of Technology, Cottbus, Germany.
[aa] Also at Forschungszentrum Jülich, Juelich, Germany.
[bb] Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland.
[cc] Also at HUN-REN ATOMKI—Institute of Nuclear Research, Debrecen, Hungary.
[dd] Also at Universitatea Babes-Bolyai—Facultatea de Fizica, Cluj-Napoca, Romania.
[ee] Also at MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary.
[ff] Also at HUN-REN Wigner Research Centre for Physics, Budapest, Hungary.
[gg] Also at Physics Department, Faculty of Science, Assiut University, Assiut, Egypt.
[hh] Also at The University of Kansas, Lawrence, Kansas, USA.
[ii] Also at Punjab Agricultural University, Ludhiana, India.
[jj] Also at University of Hyderabad, Hyderabad, India.
[kk] Also at Indian Institute of Science (IISc), Bangalore, India.
[ll] Also at University of Visva-Bharati, Santiniketan, India.
[mm] Also at IIT Bhubaneswar, Bhubaneswar, India.
[nn] Also at Institute of Physics, Bhubaneswar, India.
[oo] Also at Deutsches Elektronen-Synchrotron, Hamburg, Germany.
[pp] Also at Isfahan University of Technology, Isfahan, Iran.
[qq] Also at Sharif University of Technology, Tehran, Iran.
[rr] Also at Department of Physics, University of Science and Technology of Mazandaran, Behshahr, Iran.
[ss] Also at Department of Physics, Faculty of Science, Arak University, ARAK, Iran.
[tt] Also at Italian National Agency for New Technologies, Energy and Sustainable Economic Development, Bologna, Italy.
[uu] Also at Centro Siciliano di Fisica Nucleare e di Struttura Della Materia, Catania, Italy.
[vv] Also at Università degli Studi Guglielmo Marconi, Roma, Italy.
[ww] Also at Scuola Superiore Meridionale, Università di Napoli ’Federico II’, Napoli, Italy.
[xx] Also at Fermi National Accelerator Laboratory, Batavia, Illinois, USA.
[yy] Also at Laboratori Nazionali di Legnaro dell’INFN, Legnaro, Italy.
[zz] Also at Lulea University of Technology, Lulea, Sweden.
[aaa] Also at Consiglio Nazionale delle Ricerche—Istituto Officina dei Materiali, Perugia, Italy.
[bbb] Also at UPES—University of Petroleum and Energy Studies, Dehradun, India.
[ccc] Also at Institut de Physique des 2 Infinis de Lyon (IP2I), Villeurbanne, France.
[ddd] Also at Department of Applied Physics, Faculty of Science and Technology, Universiti Kebangsaan Malaysia, Bangi, Malaysia.
[eee] Also at Trincomalee Campus, Eastern University, Sri Lanka, Nilaveli, Sri Lanka.
[fff] Also at Saegis Campus, Nugegoda, Sri Lanka.
[ggg] Also at National and Kapodistrian University of Athens, Athens, Greece.
[hhh] Also at Ecole Polytechnique Fédérale Lausanne, Lausanne, Switzerland.
[iii] Also at Universität Zürich, Zurich, Switzerland.
[jjj] Also at Stefan Meyer Institute for Subatomic Physics, Vienna, Austria.
[kkk] Also at Near East University, Research Center of Experimental Health Science, Mersin, Turkey.
[lll] Also at Konya Technical University, Konya, Turkey.
[mmm] Also at Izmir Bakircay University, Izmir, Turkey.
[nnn] Also at Adiyaman University, Adiyaman, Turkey.
[ooo] Also at Bozok Universitetesi Rektörlügü, Yozgat, Turkey.
[ppp] Also at Istanbul Sabahattin Zaim University, Istanbul, Turkey.
[qqq] Also at Marmara University, Istanbul, Turkey.
[rrr] Also at Milli Savunma University, Istanbul, Turkey.
[sss] Also at Informatics and Information Security Research Center, Gebze/Kocaeli, Turkey.
[ttt] Also at Kafkas University, Kars, Turkey.
[uuu] Also at Istanbul Okan University, Istanbul, Turkey.
[vvv] Also at Istanbul University—Cerrahpasa, Faculty of Engineering, Istanbul, Turkey.
[www] Also at Istinye University, Istanbul, Turkey.
[xxx] Also at Yildiz Technical University, Istanbul, Turkey.
[yyy] Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom.
[zzz] Also at Monash University, Faculty of Science, Clayton, Australia.
[aaaa] Also at Università di Torino, Torino, Italy.
[bbbb] Also at Karamanoğlu Mehmetbey University, Karaman, Turkey.
[cccc] Also at California Lutheran University;, Thousand Oaks, California, USA.
[dddd] Also at California Institute of Technology, Pasadena, California, USA.
[eeee] Also at United States Naval Academy, Annapolis, Maryland, USA.
[ffff] Also at Bingol University, Bingol, Turkey.
[gggg] Also at Georgian Technical University, Tbilisi, Georgia.

[hhhh] Also at Sinop University, Sinop, Turkey.
[iiii] Also at Erciyes University, Kayseri, Turkey.
[jjjj] Also at Horia Hulubei National Institute of Physics and Nuclear Engineering (IFIN-HH), Bucharest, Romania.
[kkkk] Also at Another institute formerly covered by a cooperation agreement with CERN.
[llll] Also at Hamad Bin Khalifa University (HBKU), Doha, Qatar.
[mmmm] Also at Yerevan Physics Institute, Yerevan, Armenia.
[nnnn] Also at Imperial College, London, United Kingdom.